Vector calculus MA2VC 2013–14 — Assignment 2 SOLUTIONS

(1) We compute the total derivative of \vec{a}

$$\frac{d\mathbf{\vec{a}}}{dt}(t) = \frac{d(e^t\hat{\imath} + t\hat{\jmath} + e^{2t}\cos t\hat{k})}{dt} = e^t\hat{\imath} + \hat{\jmath} + e^{2t}(2\cos t - \sin t)\hat{k};$$

the gradient of \boldsymbol{f}

$$\vec{\nabla}f = \frac{\partial(\sqrt{x^2 + y^2 + z^2})}{\partial x}\hat{\imath} + \frac{\partial(\sqrt{x^2 + y^2 + z^2})}{\partial y}\hat{\jmath} + \frac{\partial(\sqrt{x^2 + y^2 + z^2})}{\partial z}\hat{k} = \frac{x\hat{\imath} + y\hat{\jmath} + z\hat{k}}{\sqrt{x^2 + y^2 + z^2}} = \frac{\vec{\mathbf{r}}}{|\vec{\mathbf{r}}|};$$

and their scalar product, which (thanks to the chain rule (34)) is the total derivative of the compound function:

$$\frac{d(f(\vec{a}))}{dt}(t) = \vec{\nabla}f(\vec{a}(t)) \cdot \frac{d\vec{a}}{dt}(t) \qquad \text{(chain rule)} \\
= \frac{xe^t + y + ze^{2t}(2\cos t - \sin t)}{\sqrt{x^2 + y^2 + z^2}} \\
= \frac{e^{2t} + t + e^{4t}(2\cos^2 t - \cos t\sin t)}{\sqrt{e^{2t} + t^2} + e^{4t}\cos^2 t}.$$

In the last step we have substituted the values of the curve components $a_1(t)$, $a_2(t)$ and $a_3(t)$ in place of x, y and z, so that we have a function of t only.

Alternatively, $\frac{d(f(\vec{\mathbf{a}}))}{dt}(t)$ can be found by computing and deriving the compound function:

$$\begin{aligned} f(\vec{\mathbf{a}}(t)) &= \sqrt{e^{2t} + t^2 + e^{4t}\cos^2 t},\\ \frac{d(f(\vec{\mathbf{a}}))}{dt}(t) &= \frac{2e^{2t} + 2t + 4e^{4t}\cos^2 t - 2e^{4t}\cos t\sin t}{2\sqrt{e^{2t} + t^2} + e^{4t}\cos^2 t} = \frac{e^{2t} + t + e^{4t}(2\cos^2 t - \cos t\sin t)}{\sqrt{e^{2t} + t^2} + e^{4t}\cos^2 t}. \end{aligned}$$

(2) As described in Section 2.1.1 of the notes, the length of a path can be measured by integrating the constant scalar field $f(\vec{\mathbf{r}}) = 1$. We compute the total derivative of $\vec{\mathbf{b}}$

$$\vec{\mathbf{b}} = 5\cos t\,\hat{\boldsymbol{\imath}} + (3\cos t - 4\sqrt{2}\sin t)\hat{\boldsymbol{\jmath}} + (4\cos t + 3\sqrt{2}\sin t)\hat{\boldsymbol{k}};$$
$$\frac{d\vec{\mathbf{b}}}{dt} = -5\sin t\,\hat{\boldsymbol{\imath}} + (-3\sin t - 4\sqrt{2}\cos t)\hat{\boldsymbol{\jmath}} + (-4\sin t + 3\sqrt{2}\cos t)\hat{\boldsymbol{k}};$$

its magnitude (which enters the definition of the infinitesimal length element ds)

$$\left|\frac{d\vec{\mathbf{b}}}{dt}\right|^{2} = 25\sin^{2}t + \left(-3\sin t - 4\sqrt{2}\cos t\right)^{2} + \left(-4\sin t + 3\sqrt{2}\cos t\right)^{2}$$
$$= 25\sin^{2}t + \left(9\sin^{2}t + 32\cos^{2}t + 24\sqrt{2}\sin t\cos t\right)$$
$$+ \left(16\sin^{2}t + 18\cos^{2}t - 24\sqrt{2}\sin t\cos t\right)$$
$$= 50\sin^{2}t + 50\cos^{2}t$$
$$= 50;$$

and the integral

$$\operatorname{Length}(\Gamma) = \int_{\Gamma} 1 \, \mathrm{d}t = \int_{0}^{\pi} \left| \frac{d\vec{\mathbf{b}}}{dt} \right| \, \mathrm{d}t = \int_{0}^{\pi} \sqrt{50} \, \mathrm{d}t = \pi\sqrt{50} = 5\pi\sqrt{2} \approx 22.2.$$

Since the magnitude of $\frac{d\vec{\mathbf{b}}}{dt}$ is constant $\left(\left|\frac{d\vec{\mathbf{b}}}{dt}\right| = \sqrt{50}\right)$, the integration turns out to be very easy. (The path Γ is half circle of radius $\sqrt{50}$ and $\vec{\mathbf{a}}$ travels on it with constant speed.)

(3) First, we verify that the two paths (which we denote Γ_c and Γ_d) share the endpoints:

$$\vec{\mathbf{c}}\left(-\frac{\pi}{2}\right) = \vec{\mathbf{d}}(-1) = -\hat{\imath}, \qquad \vec{\mathbf{c}}\left(\frac{\pi}{2}\right) = \vec{\mathbf{d}}(1) = \hat{\imath};$$

they both start at $-\hat{i}$ and end at \hat{i} . Then, we compute the total derivatives of the curves

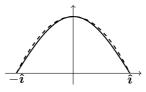
$$\frac{d\vec{\mathbf{c}}}{dt} = \frac{2}{\pi}\hat{\boldsymbol{\imath}} - \sin t\hat{\boldsymbol{\jmath}}, \qquad \qquad \frac{d\vec{\mathbf{d}}}{d\tau} = \hat{\boldsymbol{\imath}} - 2\tau\hat{\boldsymbol{\jmath}},$$

and use them to compute the integrals of $\vec{\mathbf{G}} = y\hat{\imath}$ along the two paths:

$$\begin{split} \int_{\Gamma_c} \vec{\mathbf{G}} \cdot d\vec{\mathbf{r}} &= \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \vec{\mathbf{G}} \left(\vec{\mathbf{c}}(t) \right) \cdot \frac{d\vec{\mathbf{c}}}{dt} dt \\ &= \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} (y\hat{\imath}) \cdot \left(\frac{2}{\pi} \hat{\imath} - \sin t\hat{\jmath} \right) dt = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} (\cos t\hat{\imath}) \cdot \left(\frac{2}{\pi} \hat{\imath} - \sin t\hat{\jmath} \right) dt \\ &= \frac{2}{\pi} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \cos t \, dt = \frac{2}{\pi} \sin t \Big|_{t=-\frac{\pi}{2}}^{\frac{\pi}{2}} = \frac{4}{\pi} \approx 1.273; \\ \int_{\Gamma_d} \vec{\mathbf{G}} \cdot d\vec{\mathbf{r}} &= \int_{-1}^{1} \vec{\mathbf{G}} \left(\vec{\mathbf{d}}(\tau) \right) \cdot \frac{d\vec{\mathbf{d}}}{d\tau} d\tau \\ &= \int_{-1}^{1} (y\hat{\imath}) \cdot (\hat{\imath} - 2\tau\hat{\jmath}) \, d\tau = \int_{-1}^{1} (1 - \tau^2) \hat{\imath} \cdot (\hat{\imath} - 2\tau\hat{\jmath}) \, d\tau \\ &= \int_{-1}^{1} (1 - \tau^2) \, d\tau = \left(\tau - \frac{1}{3} \tau^3 \right) \Big|_{\tau=-1}^{1} = \frac{4}{3} \approx 1.333. \end{split}$$

We note that they give different values: $\int_{\Gamma_c} \vec{\mathbf{G}} \cdot d\vec{\mathbf{r}} = \frac{4}{\pi} \neq \frac{4}{3} = \int_{\Gamma_d} \vec{\mathbf{G}} \cdot d\vec{\mathbf{r}}$. (The values are actually quite close to each other since the paths are quite similar, see Figure ??, where $\vec{\mathbf{c}}$ is plotted as a continuous line and $\vec{\mathbf{d}}$ is dashed.) Since two line integrals along different paths *sharing the endpoints* are different, by (the contrapositive of) the fundamental theorem of vector calculus (Section 2.1.3 in the notes), we have that $\vec{\mathbf{G}}$ is not a gradient, so it is not conservative.

The easiest proof of the fact that $\vec{\mathbf{G}}$ is non-conservative, consists in noting that $\vec{\nabla} \times \vec{\mathbf{G}} = -\hat{k} \neq \vec{\mathbf{0}}$, thus $\vec{\mathbf{G}}$ is not irrotational and the assertion follows from the contrapositive of the implication "conservative \Rightarrow irrotational" in the box in Section 1.5.



(4) For n = 1, 2, 3, the *n*th component of $\vec{\mathbf{H}}$ depends on the *n*th variable only, thus, by the definition of the curl operator, $\vec{\nabla} \times \vec{\mathbf{H}} = \vec{\mathbf{0}}$ and $\vec{\mathbf{H}}$ is irrotational (equivalently, we compute the curl and verify directly that $\vec{\nabla} \times \vec{\mathbf{H}} = \vec{\mathbf{0}}$).

We compute a scalar potential φ by proceeding as in Exercise 1.42:

$$\vec{\mathbf{H}} = \vec{\nabla}\varphi \quad \Rightarrow \quad \frac{\partial\varphi(x,y,z)}{\partial x} = H_1 = x \quad \Rightarrow \quad \varphi(x,y,z) = \frac{1}{2}x^2 + g(y,z),$$

$$\frac{\partial\varphi(x,y,z)}{\partial y} = \frac{\partial\left(\frac{1}{2}x^2 + g(y,z)\right)}{\partial y} = \frac{\partial g(y,z)}{\partial y} = H_2 = y^2 \quad \Rightarrow \quad g(y,z) = \frac{1}{3}y^3 + f(z),$$

$$\frac{\partial\varphi(x,y,z)}{\partial z} = \frac{\partial\left(\frac{1}{2}x^2 + \frac{1}{3}y^3 + f(z)\right)}{\partial z} = \frac{\partial f(z)}{\partial z} = H_3 = z^3 \quad \Rightarrow \quad f(z) = \frac{1}{4}z^4 + \lambda,$$

$$\Rightarrow \varphi(x,y,z) = \frac{1}{2}x^2 + g(y,z) = \frac{1}{2}x^2 + \frac{1}{3}y^3 + f(z) = \frac{1}{2}x^2 + \frac{1}{3}y^3 + \frac{1}{4}z^4 + \lambda, \quad \forall \lambda \in \mathbb{R}.$$

Both paths of Exercise (3) have endpoints $\vec{\mathbf{c}}(-\frac{\pi}{2}) = \vec{\mathbf{d}}(-1) = -\hat{\imath}$ and $\vec{\mathbf{c}}(\frac{\pi}{2}) = \vec{\mathbf{d}}(1) = \hat{\imath}$. Since $\vec{\mathbf{H}}$ is conservative, by the fundamental theorem of vector calculus, the integral along any path connecting $-\hat{\imath}$ and $\hat{\imath}$ can be computed as:

$$\int_{\Gamma_c} \vec{\mathbf{H}} \cdot d\vec{\mathbf{r}} = \int_{\Gamma_c} \vec{\nabla} \varphi \cdot d\vec{\mathbf{r}} = \int_{\Gamma_d} \vec{\mathbf{H}} \cdot d\vec{\mathbf{r}} = \int_{\Gamma_d} \vec{\nabla} \varphi \cdot d\vec{\mathbf{r}} = \varphi(\hat{\boldsymbol{\imath}}) - \varphi(-\hat{\boldsymbol{\imath}}) = \frac{1}{2} - \frac{1}{2}(-1)^2 = 0.$$