

On the long time behavior and optimal control of a tumor growth model

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BILBAO WORKSHOP ON THEORETICAL FLUID DYNAMICS

joint work with Cecilia Cavaterra (Milano), Hao Wu (Fudan)



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Outline

- 1 Phase field models for tumor growth
- 2 The model HZO by [A. Hawkins-Daarud, K.-G. van der Zee and J.-T. Oden (2011)]
- 3 Content of the joint work with C. Cavaterra and H. Wu, arXiv:1901.07500, 2019
- 4 Well-posedness
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Setting

Tumors grown *in vitro* often exhibit “layered” structures:

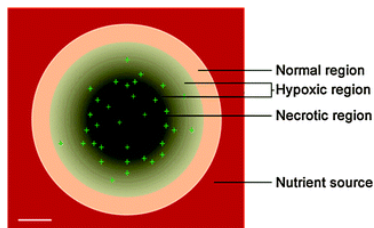


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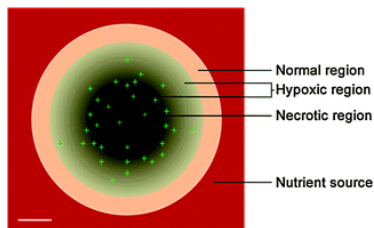


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- sharp interfaces are replaced by narrow transition layers arising due to adhesive forces among the cell species: a **diffuse interface** separates tumor and healthy cell regions
- **proliferating** tumor cells surrounded by (healthy) **host cells**, and a **nutrient** (e.g. glucose).

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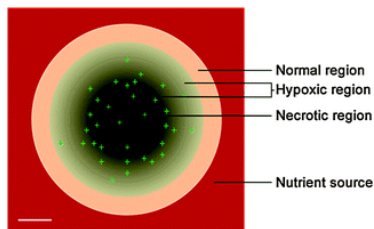


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We investigate the **two-phase** case: growth of a **tumor** in presence of a **nutrient** and surrounded by **host tissues**.

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- Ciarletta, Cristini, Frieboes, Garcke, **Hawkins-Daarud**, Hilhorst, Lam, Lowengrub, **Oden**, **van der Zee**, Wise, also for their numerical simulations → complex changes in tumor morphologies due to the interactions with nutrients or toxic agents and also due to mechanical stresses

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- Frieboes, Jin, Chuang, Wise, Lowengrub, Cristini, Garcke, Lam, Nürnberg, Sitka, for the interaction of multiple tumor cell species described by *multiphase mixture models*

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HZO: the free energy

- u = tumor cell volume fraction $u \in [0, 1]$
- n = nutrient-rich extracellular water volume fraction $n \in [0, 1]$
- $f(u) = \Gamma u^2(1 - u)^2$: a double well
- $\chi(u, n) = -\chi_0 un$: chemotaxis driving the tumor cells toward the oxygen supply

$$E = \int_{\Omega} \left(f(u) + \frac{\epsilon^2}{2} |\nabla u|^2 + \chi(u, n) + \frac{1}{2\delta} n^2 \right) dx. \quad (4)$$

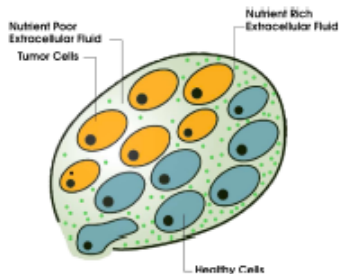


Figure 1. Four-species model: illustration of the four-species mixture. The tumor and healthy cell populations are assumed to have a thin diffuse interface, whereas the nutrient-rich and nutrient-poor extracellular water are segregated by a wide smooth interface.

The plot of the summand $f(u) + \chi(u, n)$

The lowest energy state is when $u = 1$ and $n = 1$, when there is a full interaction between the tumor species and the nutrient-rich extracellular water.

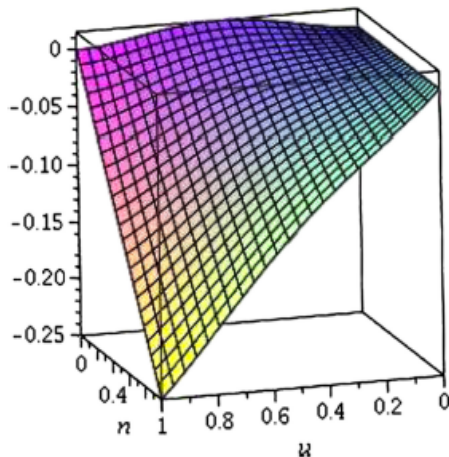


Figure 2. Graph of homogeneous free energy: $f(u) + \chi(u, n)$. ($\Gamma = \chi_0 = 0.25$).

The mass balance equations

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$$u_t = \nabla \cdot (M_u \nabla \mu_u) + \gamma_u, \quad \mu_u = \partial_u E = f'(u) + \partial_u \chi(u, n) - \epsilon \Delta u$$

$$n_t = \nabla \cdot (M_n \nabla \mu_n) + \gamma_n, \quad \mu_n = \partial_n E = \partial_n \chi(u, n) + \frac{1}{\delta} n$$

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More in particular, they choose:

$$\gamma_u = P(u)(\mu_n - \mu_u), \quad \gamma_n = -\gamma_u, \quad \text{where}$$

$$P(u) = \begin{cases} \delta P_0 u & \text{if } u \geq 0 \\ 0 & \text{elsewhere} \end{cases}$$

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Then we get

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Simulations by HZO: the tumor starts growing increasingly more ellipsoidal at first and eventually begins forming buds growing toward the higher levels of nutrient

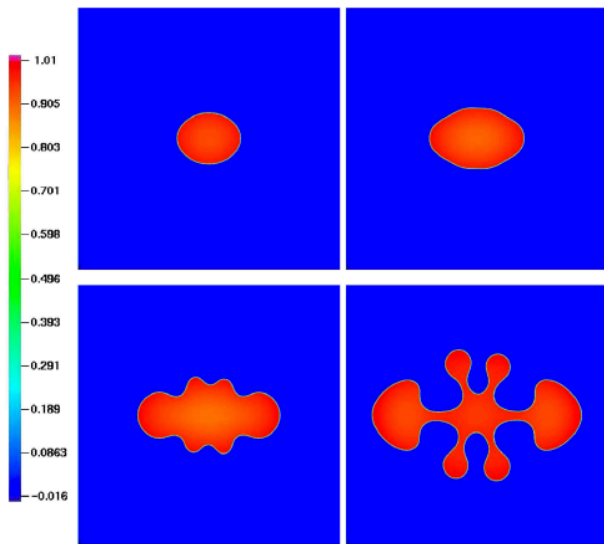


Figure 7. Example simulation: snapshots are shown at $t = 20, 40, 60$, and 80 of a simulation with $\Gamma = 0.045$, $\epsilon = 0.005$, $\chi_0 = 0.05$, $\delta = 0.01$, $P_0 = 0.1$, $\hat{M} = 200$, and $\hat{D} = 1$.

Simulations by HZO: the influence of χ_0 and δ

- When the ratio χ_0/Γ is small, the tumor remains circular $u \sim 0, 1$
- When $\chi_0 \sim \Gamma$ the tumor goes into an ellipse
- When χ_0/Γ and χ_0/ϵ are big, u no longer takes on values close to 0 and 1: it begins moving quickly toward the regions with higher nutrients
- Only when χ_0 is large the value of δ makes a difference in simulations

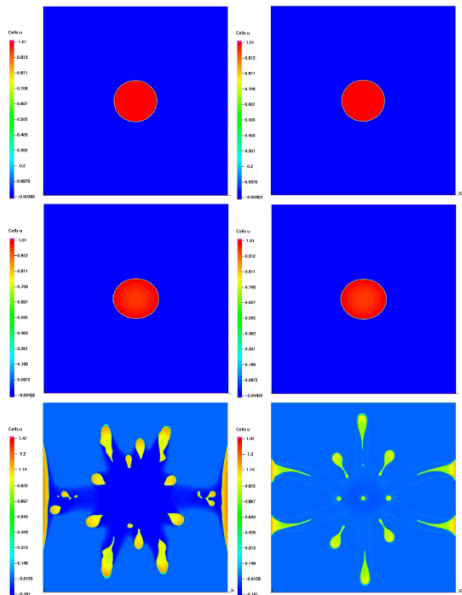
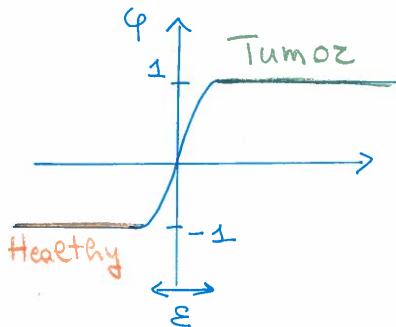


Figure 10. Effects of parameter χ_0 : illustrated here are the effects of different values of χ_0 when $\Gamma = 0.045$ and $\epsilon = 0.005$ are held constant. In the first row, $\chi_0 = 0.005$; in the second row, $\chi_0 = 0.05$; and in the third row, $\chi_0 = 0.5$. In the first column, $\delta = 0.1$; and in the second column, $\delta = 0.01$.

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Our notation for the tumor phase parameter ($u = \phi \in [-1, 1]$)



The sharp interface S replaced by a
(thickness ϵ) thin transition layer

$\phi \equiv -1$ in the Healthy tissue phase
 $\phi \equiv 1$ in the Tumor phase

Theoretical analysis: two-phase models

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- Analytical results related to well-posedness, asymptotic limits, but also **optimal control and long-time behavior of solution**, have been established in a number of papers of a number of authors which include: Agosti, Ciarletta, Colli, Frigeri, Garcke, Gilardi, Grasselli, Hilhorst, Lam, Marinoschi, Melchionna, E.R., Scala, Sprekels, Wu, etc...

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 - ▶ for tumor growth models based on the coupling of Cahn–Hilliard (for the tumor density) and reaction–diffusion (for the nutrient) equations, and
 - ▶ for models of Cahn-Hilliard-Darcy or Cahn-Hilliard-Brinkman type.

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In this talk we concentrate on two recent results on **optimal control** and **long-time behavior of solution**.

Long-time dynamics and optimal control

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By establishing the Lyapunov stability of certain equilibria of the state system (without external source), we see that ϕ_Ω can be taken as a stable configuration, so that the tumor will not grow again once the finite-time treatment is completed

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$$\begin{aligned}\phi_t - \Delta\mu &= P(\phi)(\sigma - \mu), & \mu &= -\Delta\phi + F'(\phi) \\ \sigma_t - \Delta\sigma &= -P(\phi)(\sigma - \mu) + u\end{aligned}$$

subject to initial and boundary conditions

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- The state variables are:
 - ▶ the tumor cell fraction ϕ : $\phi \simeq 1$ (tumorous phase), $\phi \simeq -1$ (healthy tissue phase)
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- F is typically a double-well potential with equal minima at $\phi = \pm 1$
- $P \geq 0$ denotes a suitable regular proliferation function
- The choice of **reactive terms** is motivated by the linear phenomenological constitutive laws for chemical reactions

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 - ▶ the nutrient concentration σ : $\sigma \simeq 1$ and $\sigma \simeq 0$ indicate a nutrient-rich or nutrient-poor extracellular water phase
- F is typically a double-well potential with equal minima at $\phi = \pm 1$
- $P \geq 0$ denotes a suitable regular proliferation function
- The choice of **reactive terms** is motivated by the linear phenomenological constitutive laws for chemical reactions
- The **control variable** u serves as an external source in the equation for σ and can be interpreted as a medication

Energy identity

The system turns out to be thermodynamically consistent. In particular, when $u = 0$ the unknown pair (ϕ, σ) is a dissipative gradient flow for the total free energy:

$$\mathcal{E}(\phi, \sigma) = \int_{\Omega} \left[\frac{1}{2} |\nabla \phi|^2 + F(\phi) \right] dx + \frac{1}{2} \int_{\Omega} \sigma^2 dx.$$

Moreover generally, under the presence of the external source u , we observe that any smooth solution (ϕ, σ) to the problem satisfies the following **energy identity**:

$$\frac{d}{dt} \mathcal{E}(\phi, \sigma) + \int_{\Omega} \left[|\nabla \mu|^2 + |\nabla \sigma|^2 + P(\phi)(\mu - \sigma)^2 \right] dx = \int_{\Omega} u \sigma dx,$$

which motivates the twofold aim of the present contribution.

Our results

1. We prove that any global weak solution will converge to a single equilibrium as $t \rightarrow +\infty$ and provide an estimate on the convergence rate.

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2. Denoting by $T \in (0, +\infty)$ a fixed maximal time in which the patient is allowed to undergo a medical treatment, we derive necessary optimality conditions for

(CP) *Minimize the cost functional*

$$\begin{aligned} \mathcal{J}(\phi, \sigma, \mathbf{u}, \tau) = & \frac{\beta_Q}{2} \int_0^\tau \int_\Omega |\phi - \phi_Q|^2 \, dx \, dt + \frac{\beta_\Omega}{2} \int_\Omega |\phi(\tau) - \phi_\Omega|^2 \, dx \\ & + \frac{\alpha_Q}{2} \int_0^\tau \int_\Omega |\sigma - \sigma_Q|^2 \, dx \, dt + \frac{\beta_S}{2} \int_\Omega (1 + \phi(\tau)) \, dx + \frac{\beta_u}{2} \int_0^\tau \int_\Omega |u|^2 \, dx \, dt + \beta_T \tau \end{aligned}$$

subject to the state system and the the control constraint

$$u \in \mathcal{U}_{\text{ad}} := \{u \in L^\infty(Q) : u_{\min} \leq u \leq u_{\max} \text{ a. e. in } Q\}, \quad \tau \in (0, T)$$

Comments on the cost functional

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- $\tau \in (0, T]$ represents the treatment time of one cycle, i.e., the amount of time the drug is applied to the patient before the period of rest, or the treatment time before surgery, ϕ_Q and σ_Q represent a desired evolution for the tumor cells and for the nutrient, ϕ_Ω stands for desired final distribution of tumor cells

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- The fifth term penalizes large concentrations of the cytotoxic drugs, and the sixth term of \mathcal{J} penalizes long treatment times

The choice of ϕ_Ω

After the treatment, the ideal situation will be either the tumor is ready for surgery or the tumor will be stable for all time without further medication (i.e., $u = 0$). This goal can be realized by making different choices of the target function ϕ_Ω in the above optimal control problem (CP).

- For the former case, one can simply take ϕ_Ω to be a configuration that is suitable for surgery.
- While for the later case, which is of more interest to us, we want to choose ϕ_Ω as a “stable” configuration of the system, so that the tumor does not grow again once the treatment is complete.

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For this purpose, we prove that any local minimizer of the total free energy \mathcal{E} is Lyapunov stable provided that $u = 0$. As a consequence, these local energy minimizers serve as possible candidates for the target function ϕ_Ω . Then after completing a successful medication, the tumor will remain close to the chosen stable configuration for all time.

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- For the single Cahn-Hilliard equation this difficulty can be overcome by employing the **Łojasiewicz-Simon** approach: a key property that plays an important role in the analysis of the Cahn-Hilliard equation is the conservation of mass, i.e.,

$$\int_{\Omega} \phi(t) \, dx = \int_{\Omega} \phi_0 \, dx \quad \text{for } t \geq 0.$$

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However, for our coupled system this property no longer holds, which brings us new difficulties in analysis.

- Besides, quite different from the Cahn-Hilliard-Oono system considered in which the mass $\int_{\Omega} \phi(t) \, dx$ is not preserved due to possible reactions, here in our case it is not obvious how to control the mass changing rate:

$$\frac{d}{dt} \int_{\Omega} \phi \, dx = \int_{\Omega} P(\phi)(\sigma - \mu) \, dx.$$

Similar problem happens to the nutrient as well, that is

$$\frac{d}{dt} \int_{\Omega} \sigma \, dx = - \int_{\Omega} P(\phi)(\sigma - \mu) \, dx + \int_{\Omega} u \, dx.$$

The problem of mass conservation

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- The observation that the **total mass** can be determined by the initial data and the external source:

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- Based on the above mentioned special structure of the system, by introducing a new version of Łojasiewicz-Simon inequality we are able to prove that every global weak solution (ϕ, σ) of the problem will converge to a certain single equilibrium $(\phi_{\infty}, \sigma_{\infty})$ as $t \rightarrow +\infty$ and, moreover, we obtain a polynomial decay of the solution.

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- Besides, a nontrivial application of the Łojasiewicz-Simon approach further leads to the Lyapunov stability of local minimizers of the free energy \mathcal{E} (we only consider the case $u = 0$ for the sake of simplicity).

Comparison with other results in the literature

- To the best of our knowledge, the only contribution in the study of long-time behavior for this problem is given in [FGR: Frigeri, Grasselli, R. (2015)] with $u = 0$, where, however, the main focus is the **existence of a global attractor** .
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Here we aim to provide a contribution to the theory of free terminal time optimal control where the control is applied in the nutrient equation.

Outline

- 1 Phase field models for tumor growth
- 2 The model HZO by [A. Hawkins-Daarud, K.-G. van der Zee and J.-T. Oden (2011)]
- 3 Content of the joint work with C. Cavaterra and H. Wu, arXiv:1901.07500, 2019
- 4 Well-posedness**
- 5 Long-term dynamics
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Let $\phi_0 \in H_N^2(\Omega) \cap H^3(\Omega)$ and $\sigma_0 \in H^1(\Omega)$ and assume that

- (P1) $P \in C^2(\mathbb{R})$ is nonnegative. There exist $\alpha_1 > 0$ and some $q \in [1, 4]$ such that, for all $s \in \mathbb{R}$, $|P'(s)| \leq \alpha_1(1 + |s|^{q-1})$
- (F1) $F = F_0 + F_1$, with $F_0, F_1 \in C^5(\mathbb{R})$. There exist $\alpha_i > 0$ and $r \in [2, 6)$ such that $|F_1''(s)| \leq \alpha_2$, $\alpha_3(1 + |s|^{r-2}) \leq F_0''(s) \leq \alpha_4(1 + |s|^{r-2})$, $F(s) \geq \alpha_5|s| - \alpha_6 \quad \forall s \in \mathbb{R}$
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Theorem (Strong solutions)

(1) For every $T > 0$, the state system admits a unique strong solution:

$$\begin{aligned} & \|\phi\|_{L^\infty(0, T; H^3(\Omega)) \cap L^2(0, T; H^4(\Omega)) \cap H^1(0, T; H^1(\Omega))} + \|\mu\|_{L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))} \\ & + \|\sigma\|_{C([0, T]; H^1(\Omega)) \cap L^2(0, T; H_N^2(\Omega)) \cap H^1(0, T; L^2(\Omega))} \leq K_1. \end{aligned}$$

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(2) Let (ϕ_i, σ_i) be two strong solutions. Then there exists a constant $K_2 > 0$, depending on $\|u_i\|_{L^2(0, T; L^2)}$, Ω , T , $\|\phi_0\|_{H^3}$ and $\|\sigma_0\|_{H^1}$, such that

$$\begin{aligned} & \|\phi_1 - \phi_2\|_{L^\infty(0, T; H^1) \cap L^2(0, T; H^3) \cap H^1(0, T; (H^1)')} + \|\mu_1 - \mu_2\|_{L^2(0, T; H^1)} \\ & + \|\sigma_1 - \sigma_2\|_{C([0, T]; H^1) \cap L^2(0, T; H^2) \cap H^1(0, T; L^2)} \leq K_2 \|u_1 - u_2\|_{L^2(0, T; L^2)}. \end{aligned}$$

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Long-term dynamics

We make the following additional assumptions:

(P2) $P(s) > 0$, for all $s \in \mathbb{R}$

(F2) $F(s)$ is real analytic on \mathbb{R}

(U2) $u \in L^1(0, +\infty; L^2(\Omega)) \cap L^2(0, +\infty; L^2(\Omega))$ and satisfies the decay condition

$$\sup_{t \geq 0} (1+t)^{3+\rho} \|u(t)\|_{L^2(\Omega)} < +\infty, \quad \text{for some } \rho > 0.$$

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Theorem (1. The stationary problem)

For any $\phi_0 \in H^1(\Omega)$, $\sigma_0 \in L^2(\Omega)$, the state system admits a unique global weak solution

(ϕ, μ, σ) : $\lim_{t \rightarrow +\infty} (\|\phi(t) - \phi_\infty\|_{H^2(\Omega)} + \|\sigma(t) - \sigma_\infty\|_{L^2(\Omega)} + \|\mu(t) - \mu_\infty\|_{L^2(\Omega)}) = 0$,

where $(\phi_\infty, \mu_\infty, \sigma_\infty)$ satisfies the stationary problem

$$\left\{ \begin{array}{ll} -\Delta \phi_\infty + F'(\phi_\infty) = \mu_\infty, & \text{in } \Omega \\ \partial_\nu \phi_\infty = 0, & \text{on } \partial\Omega \\ \int_\Omega (\phi_\infty + \sigma_\infty) dx = \int_\Omega (\phi_0 + \sigma_0) dx + \int_0^{+\infty} \int_\Omega u dx dt \end{array} \right.$$

with μ_∞ and σ_∞ being two constants given by $\sigma_\infty = \mu_\infty = |\Omega|^{-1} \int_\Omega F'(\phi_\infty) dx$.

The convergence rate

Theorem (2. Convergence rate)

Moreover, under the same assumptions, the following estimates on convergence rate hold

$$\|\phi(t) - \phi_\infty\|_{H^1(\Omega)} + \|\sigma(t) - \sigma_\infty\|_{L^2(\Omega)} \leq C(1+t)^{-\min\{\frac{\theta}{1-2\theta}, \frac{\rho}{2}\}}, \quad \forall t \geq 0,$$

$$\|\mu(t) - \mu_\infty\|_{L^2(\Omega)} \leq C(1+t)^{-\frac{1}{2} \min\{\frac{\theta}{1-2\theta}, \frac{\rho}{2}\}}, \quad \forall t \geq 0,$$

where $C > 0$ is a constant depending on $\|\phi_0\|_{H^1(\Omega)}$, $\|\sigma_0\|_{L^2(\Omega)}$, $\|\phi_\infty\|_{H^1(\Omega)}$, $\|u\|_{L^1(0,+\infty;L^2(\Omega))}$, $\|u\|_{L^2(0,+\infty;L^2(\Omega))}$ and Ω ; $\theta \in (0, \frac{1}{2})$ is a constant depending on ϕ_∞ .

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- We first derive some uniform-in-time a priori estimates on the solution (ϕ, μ, σ)
- Then we give a characterization on the ω -limit

$$\omega(\phi_0, \sigma_0) = \{(\phi_\infty, \sigma_\infty) \in (H_N^2(\Omega) \cap H^3(\Omega)) \times H^1(\Omega) : \exists \{t_n\} \nearrow +\infty \text{ such that } (\phi(t_n), \sigma(t_n)) \rightarrow (\phi_\infty, \sigma_\infty) \text{ in } H^2(\Omega) \times L^2(\Omega)\}.$$

And we have the following result

Theorem (3. The ω -limit)

Assume (P1), (F1), (U2). For any initial datum $(\phi_0, \sigma_0) \in H^1(\Omega) \times L^2(\Omega)$, the associated ω -limit set $\omega(\phi_0, \sigma_0)$ is non-empty. For any element $(\phi_\infty, \sigma_\infty) \in \omega(\phi_0, \sigma_0)$, σ_∞ is a constant and $(\phi_\infty, \sigma_\infty)$ satisfies the stationary problem. Besides, μ_∞ is a constant given by $|\Omega|^{-1} \int_\Omega F'(\phi_\infty) dx$ and the following relation holds

$$P(\phi_\infty)(\sigma_\infty - \mu_\infty) = 0, \quad \text{a.e. in } \Omega.$$

And the positivity of P entails immediately also $\sigma_\infty = \mu_\infty$.

- Finally, we prove the convergence of the trajectories and polynomial decay by means of a proper Łojasiewicz–Simon inequality:

- Finally, we prove the convergence of the trajectories and polynomial decay by means of a proper Łojasiewicz–Simon inequality: Given any initial datum $(\phi_0, \sigma_0) \in H^1(\Omega) \times L^2(\Omega)$ and source term u satisfying (U2), we denote by

$$m_\infty := |\Omega|^{-1} \left(\int_{\Omega} (\phi_0 + \sigma_0) dx + \int_0^{+\infty} \int_{\Omega} u dx dt \right)$$

the total mass at infinity time. Then we are able to derive the following

Theorem (Łojasiewicz–Simon Inequality)

Let (F1), (F2), (P1), (P2) and (U2) be satisfied. Suppose that $(\phi_\infty, \mu_\infty, \sigma_\infty)$ is a solution to the elliptic stationary problem. Then there exist constants $\theta \in (0, \frac{1}{2})$ and $\beta > 0$, depending on ϕ_∞ , m_∞ and Ω , such that for any $(\phi, \sigma) \in H_N^2(\Omega) \times H^1(\Omega)$ satisfying

$$\|\phi - \phi_\infty\|_{H^1(\Omega)} < \beta,$$

$$\int_{\Omega} (\phi + \sigma) dx + m_u |\Omega| = \int_{\Omega} (\phi_\infty + \sigma_\infty) dx = m_\infty |\Omega|,$$

where m_u is a certain constant fulfilling $|m_u| \leq |\Omega|^{-\frac{1}{2}} \|u\|_{L^1(0,+\infty;L^2(\Omega))}$, then we have

$$\begin{aligned} \|\mu - \bar{\mu}\|_{(H^1(\Omega))'} + C \|\nabla \sigma\|_{L^2(\Omega)} + C \|\sqrt{P(\phi)}(\mu - \sigma)\|_{L^2(\Omega)} + C |m_u|^{\frac{1}{2}} \\ \geq |\mathcal{E}(\phi, \sigma) - \mathcal{E}(\phi_\infty, \sigma_\infty)|^{1-\theta}, \quad \text{where} \end{aligned}$$

$\mu = -\Delta \phi + F'(\phi)$ and $C > 0$ depends on Ω , ϕ_∞ , m_∞ , $\|\phi\|_{H^2(\Omega)}$, $\|\sigma\|_{H^1(\Omega)}$, $\|u\|_{L^1(0,+\infty;L^2(\Omega))}$.

Energy minimizers with $u = 0$

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Let us now **assume** $u = 0$. Then it follows that the total mass of the system is now conserved:

$$\int_{\Omega} (\phi(t) + \sigma(t)) \, dx = \int_{\Omega} (\phi_0 + \sigma_0) \, dx, \quad \forall t \geq 0.$$

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Let $m \in \mathbb{R}$ be an arbitrary given constant. Set

$$\mathcal{Z}_m = \left\{ (\phi, \sigma) \in H^1(\Omega) \times L^2(\Omega) : \int_{\Omega} (\phi + \sigma) \, dx = |\Omega| m \right\}.$$

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Any $(\phi^*, \sigma^*) \in \mathcal{Z}_m$ is called

- a *local energy minimizer* of the total energy

$$\mathcal{E}(\phi, \sigma) = \int_{\Omega} \left[\frac{1}{2} |\nabla \phi|^2 + F(\phi) \right] \, dx + \frac{1}{2} \int_{\Omega} \sigma^2 \, dx$$

if there exists a constant $\chi > 0$ such that $\mathcal{E}(\phi^*, \sigma^*) \leq \mathcal{E}(\phi, \sigma)$, for all $(\phi, \sigma) \in \mathcal{Z}_m$ satisfying $\|(\phi - \phi^*, \sigma - \sigma^*)\|_{H^1(\Omega) \times L^2(\Omega)} < \chi$

- If $\chi = +\infty$, then (ϕ^*, σ^*) is called a *global energy minimizer* of $\mathcal{E}(\phi, \sigma)$ in \mathcal{Z}_m .

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$$\left\{ \begin{array}{ll} -\Delta\phi + F'(\phi) = \mu, & \text{in } \Omega, \\ \partial_\nu\phi = 0, & \text{on } \partial\Omega, \\ \int_\Omega (\phi + \sigma) \, dx = |\Omega|m, & \end{array} \right.$$

where μ and σ are constants given by $\sigma = \mu = |\Omega|^{-1} \int_\Omega F'(\phi) \, dx$.

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- (1) If $(\phi^*, \sigma^*) \in H_N^2(\Omega) \times \mathbb{R}$ is a strong solution to the stationary problem above, then (ϕ^*, σ^*) is a critical point of $\mathcal{E}(\phi, \sigma)$ in \mathcal{Z}_m . Conversely, if (ϕ^*, σ^*) is a critical point of $\mathcal{E}(\phi, \sigma)$ in \mathcal{Z}_m , then $\phi^* \in H_N^2(\Omega)$, $\sigma^* \in \mathbb{R}$ satisfy the stationary problem above

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- (2) If (ϕ^*, σ^*) is a local energy minimizer of $\mathcal{E}(\phi, \sigma)$ in \mathcal{Z}_m , then (ϕ^*, σ^*) is a critical point of $\mathcal{E}(\phi, \sigma)$.
- (3) The functional $\mathcal{E}(\phi, \sigma)$ has at least one minimizer $(\phi^*, \sigma^*) \in \mathcal{Z}_m$ such that

$$\mathcal{E}(\phi^*, \sigma^*) = \inf_{(\phi, \sigma) \in \mathcal{Z}_m} \mathcal{E}(\phi, \sigma)$$

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Theorem (5. Lyapunov stability)

Assume that (F1), (F2), (P1), (P2) are satisfied and $u = 0$. Given $m \in \mathbb{R}$, let (ϕ^*, σ^*) be a local energy minimizer in \mathcal{Z}_m of

$$\mathcal{E}(\phi, \sigma) = \int_{\Omega} \left[\frac{1}{2} |\nabla \phi|^2 + F(\phi) \right] dx + \frac{1}{2} \int_{\Omega} \sigma^2 dx.$$

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Then, for any $\epsilon > 0$, there exists a constant $\eta \in (0, 1)$ such that for arbitrary initial datum $(\phi_0, \sigma_0) \in (H_N^2(\Omega) \cap H^3(\Omega)) \times H^1(\Omega)$ satisfying $\int_{\Omega} (\phi_0 + \sigma_0) dx = |\Omega|m$ and $\|\phi_0 - \phi^*\|_{H^1(\Omega)} + \|\sigma_0 - \sigma^*\|_{L^2(\Omega)} \leq \eta$, the state system admits a unique global strong solution (ϕ, σ) such that

$$\|\phi(t) - \phi^*\|_{H^1(\Omega)} + \|\sigma(t) - \sigma^*\|_{L^2(\Omega)} \leq \epsilon, \quad \forall t \geq 0.$$

Namely, any local energy minimizer of $\mathcal{E}(\phi, \sigma)$ in \mathcal{Z}_m is locally Lyapunov stable.

Conclusions on long-term dynamics

- The result on long-time behavior derived in Theorem 1 and 2 can be applied to the global strong solution obtained in Theorem 5
- Although it is still not obvious to identify the asymptotic limit $(\phi_\infty, \sigma_\infty)$, we are able to conclude that $(\phi_\infty, \sigma_\infty)$ also satisfies

$$\|\phi_\infty - \phi^*\|_{H^1(\Omega)} + \|\sigma_\infty - \sigma^*\|_{L^2(\Omega)} \leq \epsilon$$

- In particular, if (ϕ^*, σ^*) is an isolated local energy minimizer then it is locally asymptotic stable

Outline

- 1 Phase field models for tumor growth
- 2 The model HZO by [A. Hawkins-Daarud, K.-G. van der Zee and J.-T. Oden (2011)]
- 3 Content of the joint work with C. Cavaterra and H. Wu, arXiv:1901.07500, 2019
- 4 Well-posedness
- 5 Long-term dynamics
- 6 The optimal control problem**
- 7 Open problems and Perspectives

Assumptions for the optimal control problem

In this section we study the optimal control problem

(CP) *Minimize the cost functional*

$$\begin{aligned} \mathcal{J}(\phi, \sigma, u, \tau) = & \frac{\beta_Q}{2} \int_0^\tau \int_\Omega |\phi - \phi_Q|^2 \, dx \, dt + \frac{\beta_\Omega}{2} \int_\Omega |\phi(\tau) - \phi_\Omega|^2 \, dx \\ & + \frac{\alpha_Q}{2} \int_0^\tau \int_\Omega |\sigma - \sigma_Q|^2 \, dx \, dt + \frac{\beta_S}{2} \int_\Omega (1 + \phi(\tau)) \, dx + \frac{\beta_u}{2} \int_0^T \int_\Omega |u|^2 \, dx \, dt + \beta_T \tau \end{aligned}$$

subject to the state system and the the control constraint

$$u \in \mathcal{U}_{\text{ad}} := \{u \in L^\infty(Q) : u_{\min} \leq u \leq u_{\max} \text{ a. e. in } Q\}, \quad \tau \in (0, T),$$

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(C1) $\beta_Q, \beta_{\Omega}, \beta_S, \beta_u, \beta_T, \alpha_Q$ are nonnegative constants but not all zero.

(C2) $\phi_Q, \sigma_Q \in L^2(Q), \phi_{\Omega}, \sigma_{\Omega} \in L^2(\Omega), u_{\min}, u_{\max} \in L^{\infty}(Q)$, and $u_{\min} \leq u_{\max}$, a.e. in Q .

(C3) Let \mathcal{U}_R be an open set in $L^2(Q)$: $\mathcal{U}_{\text{ad}} \subset \mathcal{U}_R$ and $\|u\|_{L^2(Q)} \leq R$, for all $u \in \mathcal{U}_R$.

Existence of an optimal control

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From the well-posedness results it follows that the *control-to-state operator* \mathcal{S}

$$u \mapsto \mathcal{S}(u) := (\phi, \mu, \sigma)$$

is well-defined and Lipschitz continuous as a mapping from $\mathcal{U}_R \subset L^2(Q)$ into the following space

$$(L^\infty(0, T; (H^1(\Omega))') \cap L^2(0, T; H^1(\Omega))) \times L^2(0, T; (H^1(\Omega))') \times (L^\infty(0, T; (H^1(\Omega))') \cap L^2(Q)).$$

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The triplet (ϕ, μ, σ) is the unique weak solution to the state system with data (ϕ_0, σ_0, u) over the time interval $[0, T]$. For convenience, we use the notations $\phi = \mathcal{S}_1(u)$ and $\sigma = \mathcal{S}_3(u)$ for the first and third component of $\mathcal{S}(u)$.

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The triplet (ϕ, μ, σ) is the unique weak solution to the state system with data (ϕ_0, σ_0, u) over the time interval $[0, T]$. For convenience, we use the notations $\phi = \mathcal{S}_1(u)$ and $\sigma = \mathcal{S}_3(u)$ for the first and third component of $\mathcal{S}(u)$. Then we prove the following result that implies the existence of a solution to problem (CP).

Theorem (Existence of the optimal control)

Assume that (P1), (F1), (U1) and (C1)–(C3) are satisfied. Let $\phi_0 \in H_N^2(\Omega) \cap H^3(\Omega)$ and $\sigma_0 \in H^1(\Omega)$. Then there exists at least one minimizer $(\phi_*, \sigma_*, u_*, \tau_*)$ to problem (CP).

Namely, $\phi_* = \mathcal{S}_1(u_*)$, $\sigma_* = \mathcal{S}_3(u_*)$ satisfy

$$\mathcal{J}(\phi_*, \sigma_*, u_*, \tau_*) = \inf_{\substack{(w, s) \in \mathcal{U}_{\text{ad}} \times [0, T] \\ \text{s.t. } \phi = \mathcal{S}_1(w), \sigma = \mathcal{S}_3(w)}} \mathcal{J}(\phi, \sigma, w, s).$$

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$$\begin{aligned}\partial_t \xi - \Delta \eta &= P'(\phi_*)(\sigma_* - \mu_*)\xi + P(\phi_*)(\rho - \eta), & \eta &= -\Delta \xi + F''(\phi_*)\xi, \\ \partial_t \rho - \Delta \rho &= -P'(\phi_*)(\sigma_* - \mu_*)\xi - P(\phi_*)(\rho - \eta) + h \\ \partial_n \xi &= \partial_n \eta = \partial_n \rho = 0, & \xi(0) &= \rho(0) = 0.\end{aligned}$$

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We can apply [Theorems 3.1, 3.2, CGRS] for the well-posedness of the linearized system and the Fréchet differentiability of the control-to-state operator \mathcal{S} with respect to u .

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$$\begin{aligned}\mathcal{Y} &:= \left(H^1(0, T; (H_N^2(\Omega))') \cap L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H_N^2(\Omega)) \right) \times L^2(Q) \\ &\quad \times \left(H^1(0, T; L^2(\Omega)) \cap L^2(0, T; H^2(\Omega)) \right).\end{aligned}$$

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For any $u_* \in \mathcal{U}_R$, the Fréchet derivative $D\mathcal{S}(u_*) \in \mathcal{L}(L^2(Q), \mathcal{Y})$ is defined as follows: for any $h \in L^2(Q)$, $D\mathcal{S}(u_*)h = (\xi^h, \eta^h, \rho^h)$, where (ξ^h, η^h, ρ^h) is the unique solution to the linearized system associated with h .

First order optimality conditions

Define a reduced functional

$$\tilde{\mathcal{J}}(u, \tau) := \mathcal{J}(\mathcal{S}_1(u), \mathcal{S}_3(u), u, \tau).$$

Since the control-to-state mapping \mathcal{S} is also Fréchet differentiable into $C^0([0, T]; L^2(\Omega))$ with respect to u , then the reduced cost functional $\tilde{\mathcal{J}}$ is Fréchet differentiable in \mathcal{U}_R .

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Theorem (Existence of solutions to the adjoint system)

Assume (P1), (F1), (U1), (C1)–(C3), $\phi_0 \in H_N^2(\Omega) \cap H^3(\Omega)$, and $\sigma_0 \in H^1(\Omega)$. Then the adjoint system

$$\begin{aligned} -\partial_t p + \Delta q - F''(\phi_*) q + P'(\phi_*)(\sigma_* - \mu_*)(r - p) &= \beta_Q (\phi_* - \phi_Q) \\ q - \Delta p + P(\phi_*)(p - r) = 0, \quad -\partial_t r - \Delta r + P(\phi_*)(r - p) &= \alpha_Q (\sigma_* - \sigma_Q) \\ \partial_n p = \partial_n q = \partial_n r = 0, \quad r(\tau_*) = 0, \quad p(\tau_*) &= \beta_\Omega (\phi_*(\tau_*) - \phi_\Omega) + \frac{\beta_S}{2} \end{aligned}$$

has a unique weak solution (p, q, r) on $[0, \tau_*]$:

$$\begin{aligned} p &\in H^1(0, \tau_*; (H_N^2(\Omega))') \cap C^0([0, \tau_*]; L^2(\Omega)) \cap L^2(0, \tau_*; H_N^2(\Omega)), \\ q &\in L^2(\Omega \times (0, \tau_*)), \quad r \in H^1(0, \tau_*; L^2(\Omega)) \cap C^0([0, \tau_*]; H^1(\Omega)) \cap L^2(0, \tau_*; H_N^2(\Omega)). \end{aligned}$$

Necessary optimality conditions

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Theorem (Necessary optimality conditions)

Let $(u_*, \tau_*) \in \mathcal{U}_{\text{ad}} \times [0, T]$ denote a minimizer to the optimal control problem (CP) with corresponding state variables $(\phi_*, \mu_*, \sigma_*) = \mathcal{S}(u_*)$ and associated adjoint variables (p, q, r) , then it holds:

$$\beta_u \int_0^T \int_{\Omega} u_*(u - u_*) \, dx \, dt + \int_0^{\tau_*} \int_{\Omega} r(u - u_*) \, dx \, dt \geq 0, \quad \forall u \in \mathcal{U}_{\text{ad}}.$$

Besides, setting

$$\begin{aligned} \mathcal{L}(\phi_*, \sigma_*, \tau_*) &= \frac{\beta_Q}{2} \int_{\Omega} |\phi_*(\tau_*) - \phi_Q(\tau_*)|^2 \, dx + \beta_{\Omega} \int_{\Omega} (\phi_*(\tau_*) - \phi_{\Omega}) \partial_t \phi_*(\tau_*) \, dx \\ &+ \frac{\alpha_Q}{2} \int_{\Omega} |\sigma_*(\tau_*) - \sigma_Q(\tau_*)|^2 \, dx + \frac{\beta_S}{2} \int_{\Omega} \partial_t \phi_*(\tau_*) \, dx + \beta_T \end{aligned}$$

we have

$$\mathcal{L}(\phi_*, \sigma_*, \tau_*) \begin{cases} \geq 0, & \text{if } \tau_* = 0, \\ = 0, & \text{if } \tau_* \in (0, T), \\ \leq 0, & \text{if } \tau_* = T. \end{cases}$$

Interpretation of the first condition

Besides, if we extend r by zero to $(\tau_*, T]$, then we can express the variational inequality

$$\beta_u \int_0^T \int_{\Omega} u_*(u - u_*) \, dx \, dt + \int_0^{\tau_*} \int_{\Omega} r(u - u_*) \, dx \, dt \geq 0, \quad \forall u \in \mathcal{U}_{\text{ad}}.$$

as

$$\int_0^T \int_{\Omega} (\beta_u u_* + r)(u - u_*) \, dx \, dt \geq 0, \quad \forall u \in \mathcal{U}_{\text{ad}},$$

which allows the interpretation that the optimal control u_* is the $L^2(Q)$ -projection of $-\beta_u^{-1}r$ onto the set \mathcal{U}_{ad} (provided that $\beta_u > 0$).

Outline

- 1 Phase field models for tumor growth
- 2 The model HZO by [A. Hawkins-Daarud, K.-G. van der Zee and J.-T. Oden (2011)]
- 3 Content of the joint work with C. Cavaterra and H. Wu, arXiv:1901.07500, 2019
- 4 Well-posedness
- 5 Long-term dynamics
- 6 The optimal control problem
- 7 Open problems and Perspectives

Open problems and Perspectives

- O1. In practice it would be safer for the patient (and thus more desirable) to approximate the target functions in the L^∞ -sense rather than in the L^2 -sense or to include a **pointwise state constraint** on ϕ : $|\phi(x, \tau) - \phi_\Omega| \leq \epsilon$ for a.e. $x \in \Omega$. This leads to a more involved adjoint system.

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- P4. **Include a stochastic** term in phase-field models for tumor growth representing for example uncertainty of a therapy or random oscillations of the tumor phase (with C. Orrieri and L. Scarpa).

Many thanks to all of you for the attention!

- Def. \mathcal{B}_0 is an *absorbing set* for a semigroup $S(t)$ on a metric space (X, d_X) iff
 - ▶ \mathcal{B}_0 is bdd
 - ▶ $\forall B \subset X$ bdd $\exists T_B \geq 0$ s.t. $S(t)B \subset \mathcal{B}_0 \quad \forall t \geq T_B$.
- Theorem. Let $S(t)$ be a strongly continuous semigroup on a c.m.s. (X, d_X) . Moreover, if
 - ▶ $S(t)$ admits an absorbing set \mathcal{B}_0 ;
 - ▶ $\forall B \subset X$ bdd $\exists t_B > 0$ s.t. $\bigcup_{t \geq t_B} S(t)B$ is compact in X ,then $S(t)$ admits a *universal attractor* \mathcal{A} that is

$$\mathcal{A} = \bigcap_{\tau \geq 0} \overline{\bigcup_{t \geq \tau} S(t)\mathcal{B}_0}.$$